# High-speed pulse train generation by spectral filtering of a mode-locked laser output

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#### Abstract

We demonstrate a simple method for generating high-speed pulse trains from a relatively low-speed pulse source by spectral filtering using a high-finesse fibre Fabry-Pérot interferometer. Reasonably stable 40 GHz and 100 GHz pulse trains are obtained at 1550 nm from a 10 GHz mode-locked fibre ring laser.

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### 1. Introduction

High-speed optical pulses are very attractive for many applications, such as microwave/millimeter-wave photonics, high-speed optical-time-division-multiplexing optical communications, high-speed optical sampling, etc. Active mode locking of fibre ring lasers is an efficient technique for generating stable picosecond pulses with low-timing jitter at several GHz. However, the maximum pulse speed of actively mode-locked fibre lasers (MLFLs) is practically limited by the bandwidth of modulators and/or their driving electronics. It is therefore a challenging problem to find ways for increasing repetition speed of the output pulses from an actively MLFL, without increasing the intensity or phase modulation frequency.

Several techniques for increasing the pulse speed by repetition rate multiplication have already been suggested. A rational harmonic mode-locking technique has offered successful generation of optical pulses at several times of the driving frequency of intensity modulator [1, 2]. However, the output has shown serious periodic pulse amplitude fluctuations with increasing multiplication factor. Intra-cavity filtering or higher-order frequency modulation mode-locking based on high-finesse Fabry-Pérot interferometers has resulted in a good stability [3, 4]. However, the multiplication factor has been relatively small and the inserted interferometer has induced large laser cavity loss.

Without changing anything inside the laser cavity, high-speed pulse trains could also be generated by simply passing a mode-locked laser output through a device or apparatus for repetition rate multiplication, such as a dispersive fibre [5], fibre Bragg grating [6], a time-to-space pulse shaper [7], etc. Those methods also suffer from some drawbacks, such as pulse-to-pulse spectrum change [4, 5] and/or relatively large amplitude variation [6–8]. An arrayed waveguide grating with a relatively narrow free spectral range (FSR) has been used as an extra-cavity repetition rate multiplier by periodic spectral filtering of ultrashort pulses in its FSR [8]. Physically small FSR implemented by large (larger than the input pulse width) time-delay increment per waveguide split individual input pulses into multiple pulses to achieve repetition rate multiplication corresponding to the number of waveguides. Even though successful repletion rate multiplication has been achieved in the work [8], each pass-band spectrum has shown side bands, due to a wide bandwidth of the arrayed waveguide grating filter, resulting in pulse burst (rather than continuous pulse train) with the period of input pulses.

Repetition rate multiplication by double-passing a low-finesse Fabry-Pérot interferometer has also been reported [9]. The amplitude fluctuation in the output pulse train has been reduced at the cost of extra components (Faraday rotator, polarising beam splitter, polarisation controller, etc.) for doubly passing the Fabry-Pérot interferometer. In the present study we use a high-finesse fibre Fabry-Pérot interferometer (FFPI) for a single pass but similar sharp spectral filtering, ensuring stable pulse train generation within its FSR. As a result, we obtain reasonably stable 40 GHz and 100 GHz pulse trains from a 10 GHz mode-locked fibre.

#### 2. Experimental setup

Fig. 1 shows a conceptual diagram of high-speed pulse train generation from a relatively slow pulse source by extra-cavity spectral filtering. A mode-locked laser output at a pulse repetition frequency f (the period T = 1/f) reveals a discrete spectral comb with the frequency spacing f, as shown in Fig. 1. Using an appropriate filter, such as a Fabry-Pérot etalon, we can select periodic spectral lines with the frequency spacing mf, where m is a positive integer. Then the resultant pulse speed increases m times (mf) when compare with the original speed, while the pulse period decreases to T/m (see Fig. 1). In this way we can get a high-speed pulse train from a relatively slow pulse train by a simple spectral filtering.



**Fig. 1.** Schematic diagram of high-speed pulse train generation by spectral filtering: spacing of the spectral comb increases by *m* times when compare with the original spacing *f*, resulting in increasing pulse speed (*mf*).



Note that the envelope of the spectral comb does not change after the spectral filtering, ensuring the same pulse shape and width after multiplying *m* times the pulse repetition rate.

с

1552.5

Fig. 2. Schematic diagram of our experimental setup.

Fig. 2 shows schematic experimental setup confirming our idea. A mode-locked fibre ring laser (MLFL) generates stable 10 GHz pulse trains with 2 ps pulse width at 1550 nm. For periodic spectral filtering, a MLFL output passes through a high-finesse (~ 33) FFPI with the 40 GHz or 100 GHz FSR. After passing the FFPI, the output is amplified by an erbium-doped fibre amplifier and its characteristics are measured by an optical spectrum analyser, a radio frequency (RF) spectrum analyser coupled with a fast photodetector, an autocorrelator, and a fast sampling oscilloscope. The MLFL pulse frequency is tuned deliberately to match its harmonics with the FSR of the FFPI.

## **3. Experimental results**

As discussed above, a 10 GHz MLFL output is characterised by evenly spaced discrete spectral lines (an optical-frequency comb), with the frequency spacing equal to the pulse



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repetition rate at 10 GHz. After passing the high-finesse FFPI with 40 or 100 GHz FSR, we can filter out any unwanted spectral components from the 10 GHz spectral comb, generating a new spectral comb with the 40 GHz or 100 GHz spacing, as demonstrated in Fig. 3b or Fig. 3c. As expected, the spectral envelopes of the 40 and 100 GHz combs have their shapes very similar with that of 10 GHz, thus ensuring the same pulse shape/width after the spectral filtering, *i.e.* after the repetition rate multiplication.



To double confirm the fact of high-speed pulse train generation, we have measured the auto/cross correlation traces of the filtered outputs by delaying one of the correlation signals, as shown in Fig. 4. The results of those measurements substantiate that the pulse period is reduced from 100 ps to 25 ps or 10 ps. In addition, they also show the input and output pulse widths to be almost the same. Again we make sure that the 40 GHz and 100 GHz pulse trains are generated from the 10 GHz MLFL by simple extra-cavity spectral filtering.

In order to check stability of the generated pulse train in details, we have detected the optical signal using a fast detector and examined its RF spectrum (see Fig. 5). For comparison, we have kept the same optical power of 1.5 dBm at the RF spectrum analyser when measuring the RF spectra of the input and FPPI output pulse trains. Here the RF spectra are measured up to 25 GHz, the limit of our measurement facility. The 10 GHz input pulses show clear 10 GHz and its second harmonic components, as shown in Fig. 5a.

Note that 10 GHz and its harmonics (see Fig. 5b) are completely suppressed in the generated 40 GHz pulses, showing clean enough pulses. However, a noticeable 10 GHz component is observed in the 100 GHz train, even though it is suppressed more than 20 dB, as depicted in Fig. 5c. We think that the 10 GHz component, indicating its periodic amplitude modulation at 10 GHz, is induced from the limited finesse of the 100 GHz FFPI. This is understandable since the FFPI with broader FSR (for larger number of repetition rate multiplication) requires higher finesse to provide similar output quality. Similar envelope modulation, due to a finite finesse (or broad pass-band width) of a spectral filter, has been observed in the arrayed waveguide grating filter reported in [8].





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We have also measured directly the time traces of the 10 GHz input and the 40 GHz FFPI output, using a 40 GHz fast sampling scope (see Fig. 6). Again, reasonably stable 40 GHz signal is observed. We believe that a small 10 GHz amplitude fluctuation of the 40 GHz train, which appears to be completely suppressed in the RF spectrum measurement (see Fig. 5b), is produced by some unknown 10 GHz system operation noise and/or a limited FFPI finesse.

## 4. Conclusion

In this paper we have suggested a simple technique aimed at generation of high-speed pulse trains from relatively slow pulse sources, using spectral filtering with a high-finesse Fabry-Pérot interferometer for repetition rate multiplication. The FFPI selects a spectral comb with the frequency spacing corresponding to its FSR to generate a high-speed pulse train within its FSR. The optical spectra measured, autocorrelation traces, the RF spectra, and the sampling scope traces have revealed reasonably stable 40 GHz and 100 GHz pulse train generation from the 10 GHz mode-locked laser by the spectral filtering technique suggested for repetition rate multiplication. The fundamental 10 GHz component is suppressed completely in the 40 GHz train and more than 20 dB in the 100 GHz train. Other arbitrary-speed pulse trains at the harmonics of the input signal can be easily generated with the similar spectral filtering technique.

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Анотація. Запропоновано простий метод генерації високошвидкісних серій імпульсів від відносно низькошвидкісного імпульсного джерела шляхом спектрального фільтрування з використанням високодобротного волоконного інтерферометра Фабрі-Перо. Отримано стабільні серії імпульсів 40 ГГц і 100 ГГц з довжиною хвилі 1550 нм від 10 ГГц-волоконного кільцевого лазера з синхронізацією мод.